

Momentum conservation and collisions

Impulse

- 1 Incident momentum is 12 kg m s^{-1} . Rebound momentum is -8 kg m s^{-1} . Impulse on ball is -20 kg m s^{-1} . Impulse on post is $+20 \text{ kg m s}^{-1}$.
- 2 (a) Incident momentum is $1.40 \times 10^{-23} \text{ kg m s}^{-1}$. Rebound momentum is $-1.40 \times 10^{-23} \text{ kg m s}^{-1}$. Impulse on wall is $+2.79 \times 10^{-23} \text{ kg m s}^{-1}$. Force on wall is
(b) Force on wall is 27.9 kN .

Conservation of momentum and energy.

- 3 Consider a small time interval dt at the start of which the rocket has mass m and velocity v and at the end of which it has mass $m + dm$ and velocity $v + dv$. The products of mass $-dm$ are ejected with a velocity $v - v_r$. (Note dm is negative). By conservation of momentum $mv = (m + dm)(v + dv) - dm(v - v_r)$, which simplifies to first order $0 = mdv + v_r dm$.

$$\text{Hence } \int_{v_1}^{v_2} \frac{dv}{v_r} = - \int_{m_1}^{m_2} \frac{dm}{m}, \text{ or } v_2 = v_1 + v_r \ln \frac{m_1}{m_2}.$$

- 4 (a) $\text{KE} = 2.09 \times 10^{-20} \text{ J}$, and photon energy is $5.68 \times 10^{-19} \text{ J}$
(b) Momentum is conserved, the initial momentum is $2m_{Cl}v_x \mathbf{i}$, one atom is scattered with momentum $m_{Cl}v_y \mathbf{j}$, thus $2m_{Cl}v_x \mathbf{i} = m_{Cl}v_y \mathbf{j} + \mathbf{p}$ so that $\mathbf{p} = 2m_{Cl}v_x \mathbf{i} - m_{Cl}v_y \mathbf{j}$, and $\mathbf{v} = 2v_x \mathbf{i} - v_y \mathbf{j} = (1200, -1600) \text{ m s}^{-1}$
(c) The atoms have $\text{KE} 7.43 \times 10^{-20} \text{ J}$ and $1.16 \times 10^{-19} \text{ J}$, the total is $1.90 \times 10^{-19} \text{ J}$.
(d) Binding energy for the bond that has been broken.
- 5 Momentum is conserved, hence $2mv_f = mv$, so $v_f = v/2$

- 6 Momentum is conserved, and is equal to $(7.37, 6.90) \times 10^{-23} \text{ kg m s}^{-1}$. The velocity is obtained by dividing this vector by the mass of the complex, and is $(249, 234) \text{ m s}^{-1}$. (The speed is 342 m s^{-1} .)

Elastic collisions.

- 7 (a) A collision in which both momentum and kinetic energy are conserved.

(b) (i) Denote the incident velocity of the neutron by \mathbf{u} , before the collision the momentum of the neutron plus the C atom is $m_n \mathbf{u}$ because the C atom is stationary, and the velocity of the neutron relative to the C atom is \mathbf{u} . After the collision the total momentum is unchanged, i.e. $m_n \mathbf{v}_n + m_c \mathbf{v}_c = m_n \mathbf{u}$ and the relative velocity will have changed sign, $\mathbf{v}_n - \mathbf{v}_c = -\mathbf{u}$. Solving,

$$\mathbf{v}_n = -\frac{(m_c - m_n)}{(m_n + m_c)} \mathbf{u} = -2.2 \times 10^7 \text{ m s}^{-1} \text{ and } \mathbf{v}_c = \frac{2m_n}{(m_c + m_n)} \mathbf{u} = 4 \times 10^6 \text{ m s}^{-1}, \text{ So the}$$

(ii) The initial and final energies of the neutron are $5.6 \times 10^{-13} \text{ J}$ and $4.0 \times 10^{-13} \text{ J}$, so the energy lost is $1.6 \times 10^{-13} \text{ J}$ (28% of the total).

(c) The maximum energy loss will be all the energy of the neutron. This will only be possible if $\frac{(m_x - m_n)}{(m_n + m_x)} = 0$, i.e. the nucleus should have a mass of 1 unit (which is why H atoms in water are widely used for this purpose).

- 8 (a) The balls have the same mass, momentum is conserved. Let the initial velocity of the cue ball be \mathbf{u} in the x -direction, and its scattering speed and angle be v_c and θ , respectively. And let the scattering speed and angle of the object ball be v_o and ϕ , respectively.

Before the collision the total momentum vector is $\mathbf{p} = m\mathbf{u}(1, 0)$ and the relative velocity is (cue – object) $\mathbf{u}_{rel} = u(1, 0)$.

After the collision the momentum is unchanged and the relative velocity has the same magnitude but a different direction. Hence we get the equations

$$\mathbf{p} = m\mathbf{u}(1, 0) = mv_c(\cos\theta, \sin\theta) + mv_o(\cos\phi, \sin\phi) \text{ and}$$

$$|\mathbf{v}_{rel}| = |v_c(\cos\theta, \sin\theta) - v_o(\cos\phi, \sin\phi)| = |(v_c \cos\theta - v_o \cos\phi, v_c \sin\theta - v_o \sin\phi)| = |\mathbf{u}_{rel}| = u$$

Conservation of momentum gives

$$v_c \cos\theta + v_o \cos\phi = u$$

$$v_c \sin\theta + v_o \sin\phi = 0$$

which can be used to substitute for the properties of the object ball giving,

$$|(2v_c \cos \theta - u, 2v_c \sin \theta)| = u \Rightarrow v_c = u \cos \theta. \text{ And now that } v_c \text{ is known } v_o \text{ can be found}$$

$$\begin{aligned} v_o \cos \phi &= u - v_c \cos \theta \\ v_o \sin \phi &= -v_c \sin \theta \end{aligned} \quad \text{so that } v_o^2 = u^2 - u^2 \cos^2 \theta, \text{ or equivalently } v_o = u \sin \theta,$$

$$\text{and finally } \tan \phi = -\cot \theta.$$

Putting in the numbers the cue ball is scattered with a speed of 4 m s^{-1} at an angle of 37° , and the object ball with a speed of 3.4 m s^{-1} at an angle of -53° . Note that the balls are moving in perpendicular directions. This is implied by the relationship between the two angles.

(b) Before the collision the total momentum vector is $\mathbf{p} = \frac{1}{2} m u (1, 0)$ and the relative velocity is (cue - object) $\mathbf{u}_{rel} = u(1, 0)$.

Hence we get the equations $\mathbf{p} = \frac{1}{2} m u (1, 0) = \frac{1}{2} m v_c (\cos \theta, \sin \theta) + m v_o (\cos \phi, \sin \phi)$ and

$$|\mathbf{v}_{rel}| = |v_c (\cos \theta, \sin \theta) - v_o (\cos \phi, \sin \phi)| = |(v_c \cos \theta - v_o \cos \phi, v_c \sin \theta - v_o \sin \phi)| = |\mathbf{u}_{rel}| = u$$

Conservation of momentum gives

$$\frac{1}{2} v_c \cos \theta + v_o \cos \phi = \frac{1}{2} u$$

$$\frac{1}{2} v_c \sin \theta + v_o \sin \phi = 0$$

which can be used to substitute for the properties of the object ball giving,

$$|\left(\frac{3}{2} v_c \cos \theta - \frac{1}{2} u, v_c \frac{3}{2} v_c \sin \theta\right)| = u \Rightarrow 3v_c^2 - 2uv_c \cos \theta - u^2 = 0, \text{ a quadratic whose solution is}$$

$$v_c = \frac{1}{3} u \left(\cos \theta + \sqrt{3 + \cos^2 \theta} \right) = 4.51 \text{ m s}^{-1}.$$

And now that v_c is known v_o can be found

$$\begin{aligned} 2v_o \cos \phi &= u - v_c \cos \theta \\ 2v_o \sin \phi &= -v_c \sin \theta \end{aligned} \quad \text{so that } v_o^2 = \frac{1}{4} (u^2 - 2uv_c \cos \theta + v_c^2) = 1.53 \text{ m s}^{-1},$$

$$\text{and finally } \tan \phi = \frac{-v_c \sin \theta}{u - v_c \cos \theta} \Rightarrow \theta = -63^\circ.$$

- 9 At the closest approach the protons turn round, and so are momentarily stationary. All their KE has been converted to PE of Coulombic repulsion, i.e.

$$m_p v^2 = \frac{e^2}{4\pi\epsilon_0 r} \Rightarrow r = \frac{e^2}{4\pi\epsilon_0 m_p v^2} = 3.47 \text{ } \mu\text{m}$$

It is better to use the relative motion to analyse this problem, because it is easy to apply to the

case where the speeds are not equal. $\frac{1}{2} \mu v_{rel}^2 = \frac{e^2}{4\pi\epsilon_0 r}$. The result is the same.

- 10 (a) At closest approach the relative KE is zero, so the initial relative KE is all converted to PE.

$$\frac{1}{2}\mu u_{rel}^2 = 4\varepsilon \left(\left(\frac{\sigma}{r_0} \right)^{12} - \left(\frac{\sigma}{r_0} \right)^6 \right) = 4\varepsilon (x^2 - x) \Rightarrow x^2 - x - \frac{\mu u_{rel}^2}{8\varepsilon} = 0 \Rightarrow x = \frac{1}{2} \left(1 + \sqrt{1 + \frac{\mu u_{rel}^2}{2\varepsilon}} \right) = 1.3$$

hence $r/\sigma = 0.957$ and $r = 3.25 \text{ \AA}$.

- (b) At closest approach the relative velocity is zero so both atoms are travelling with the same speed, the centre of mass speed, 200 m s^{-1} .

- (c) The repulsive force is the derivative of the PE,

$$F = -\frac{dV}{dr} = \frac{4\varepsilon}{r} \left(12 \left(\frac{\sigma}{r} \right)^{12} - 6 \left(\frac{\sigma}{r} \right)^6 \right) = 260 \text{ pN}, \text{ thus the relative acceleration is}$$

$F/\mu = 7.88 \times 10^{15} \text{ m s}^{-2}$. The accelerations of each atom are equal and opposite and half of this.

- (d) The momentum is conserved and the final relative velocity will be equal and opposite to the initial relative velocity. Thus

$$u_A + u_B = v_A + v_B \text{ and } u_A - u_B = -v_A + v_B. \text{ Hence } v_A + v_B = 400 \text{ m s}^{-1} \text{ and}$$

$-v_A + v_B = 400 \text{ m s}^{-1}$, solving, $v_B = 400 \text{ m s}^{-1}$ and $v_A = 0 \text{ m s}^{-1}$. The incident particle transfers all its energy to the target.

Kinetic theory of pressure and effusion (application of mechanics).

- 11 (a) 297 m s^{-1} . (Common errors are to use g instead of kg for mass, and to forget that nitrogen is diatomic).

(b) In time dt all molecules in the cylinder of length $v dt$ and cross sectional area A travelling towards the wall will hit it, the volume of the cylinder is $Av dt$, and it therefore contains $NAv dt/2$ particles (half molecules in the volume travelling in each direction). Thus in time t there will be $NAvt/2 = NAt\sqrt{RT/M}/2$ collisions.

(c) Each molecule changes momentum from mv to $-mv$. The total rate of change of momentum of gas molecules is therefore $-mNART/M$, and the force on the wall will be equal and opposite, $F = mNART/M$, so the pressure is $p = mNRT/M = NRT/N_A = nRT/V$

(d) By symmetry, gases are isotropic, so the mean square velocity component in any direction must be the same.

$$p = \frac{1}{3} N m \overline{c^2}$$

12 (a) $Nf(v_x)dv_x$

(b) $Nf(v_x)dv_x \times Av_x dt$

(c) Rate of change of molecular momentum for this subgroup is

$$Nf(v_x)dv_x \times Av_x \times (-2mv_x) = -2NmAv_x^2 f(v_x)dv_x, \text{ hence contribution to force is}$$

$$2NmAv_x^2 f(v_x)dv_x$$

(d) $F = 2NA \int_0^{\infty} v_x^2 f(v_x) dv_x$. But because the pdf is an even function of v_x and the integrals over

negative and positive v_x are equal, hence $F = NmA \int_{-\infty}^{\infty} v_x^2 f(v_x) dv_x = NmA \overline{v_x^2}$.

A more sophisticated version is possible using the full 3d probability density for the velocity.